

Canonical forms of two-qubit states under local operations

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Abstract: We provide a comprehensive analysis to obtain algebraically distinct canonical forms of two-qubit density matrices under local operations. We identify two inequivalent canonical forms realized through Lorentz singular value decompositions of the real parametrization of two-qubit density matrix. An elegant geometric visualization of two-qubit states on and within the Bloch ball follows naturally from our approach.

Keywords: Local operations, Canonical forms, Lorentz invariance, Bloch ball

The action $\rho_{AB} \rightarrow A \otimes B \rho_{AB} A^\dagger \otimes B^\dagger$ (upto a normalization factor) of stochastic local operation and classical communication (SLOCC) on a two-qubit density matrix $\rho_{AB} = \frac{1}{4} \sum_{\mu, \nu=0}^3 \Lambda_{\mu\nu} (\sigma_\mu \otimes \sigma_\nu)$ manifests as Lorentz transformation on the 4×4 real matrix Λ (defined by $\Lambda_{\mu\nu} = \text{Tr} [\rho_{AB} (\sigma_\mu \otimes \sigma_\nu)]$) in the form $\Lambda \rightarrow \bar{\Lambda} = L_A \Lambda L_B^T$, where $A, B \in \text{SL}(2, \mathbb{C})$ and $L_A, L_B \in \text{SO}(3, 1)$. With the help of suitable Lorentz transformations L_{A_c}, L_{B_c} , it is possible to arrive the Lorentz singular value decomposition $\Lambda^c = L_{A_c} \Lambda L_{B_c}^T$ of the real matrix Λ parametrizing the two-qubit system. Verstraete et.al. [1, 2] had arrived at two different types of Lorentz singular value decomposition for the real matrix Λ under SLOCC, employing highly technical results on matrix decompositions in n dimensional spaces with indefinite metric [3]. It was pointed out that the canonical forms of Refs.[1, 2] fail to reveal the underlying geometric features in an unambiguous way [4]. In this work we have carried out a complete analysis for obtaining the algebraically different canonical forms of two-qubit states based on an entirely different approach – inspired by the techniques developed in classical polarization optics by some of us [5, 6]. Our approach leads to an elegant geometrical representation of two-qubit state on and within the Bloch ball [7].

We construct two real symmetric matrices $\Omega_A = \Lambda G \Lambda^T$, and $\Omega_B = \Lambda^T G \Lambda$ where $G = \text{diag}(1, -1, -1, -1)$ denotes the Minkowski metric. These matrices undergo Lorentz congruent transformations $\Omega_{A,B} \rightarrow \bar{\Omega}_{A,B} = L_{A,B} \Omega_{A,B} L_{A,B}^T$. Following the detailed mathematical analysis carried out in Refs. [5, 6], we recognize that the matrices $G \Omega_{A,B}$ [7] are *positive semidefinite* and they have *either time-like or light-like Minkowski eigenvectors associated with their highest eigenvalue*. This results in *two* distinct canonical forms Λ^{lc} or Λ^{llc} for the real matrix Λ . The type-I canoni-

cal form, corresponding to time-like eigenvector associated with the highest eigenvalue λ_0 of $G \Omega_{A,B}$, is diagonal $\Lambda^{lc} = \text{diag}(1, \sqrt{\lambda_1/\lambda_0}, \sqrt{\lambda_2/\lambda_0}, \sqrt{\lambda_3/\lambda_0})$ and the associated two-qubit state is in the Bell-diagonal form [7]. On the other hand the type-II canonical form Λ_A^{llc} (corresponding to light-like eigenvector for the highest eigenvalue of $G \Omega_{A,B}$ has a non-diagonal form given by [7] $\Lambda_A^{llc} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & r_1 & 0 & 0 \\ 0 & 0 & -r_1 & 0 \\ 1-r_0 & 0 & 0 & r_0 \end{pmatrix}$, where $0 \leq r_1^2 \leq r_0 \leq 1$, $r_\alpha = \frac{\lambda_\alpha}{\lambda_0}$, $\mu = 0, 1$, $\phi_0 = (L^T A \Omega_A L_A)_{00}$, with $\lambda_0 \geq \lambda_1$ denoting doubly degenerate eigenvalues of $G \Omega_A$. We show that under the map Λ^{lc}

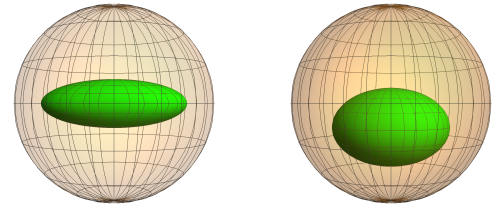


FIG. 1: Ellipsoidal (left) and spherodal (right) surfaces inside the Bloch ball characterizing type-I, II canonical forms $\Lambda_A^{lc}, \Lambda_A^{llc}$ parametrizing the SLOCC inequivalent families of two-qubit density matrices.

the set of all four-vectors $\{(1, x_1, x_2, x_3)^T | x_1^2 + x_2^2 + x_3^2 = 1\}$ representing the Bloch sphere gets transformed to $\{\Lambda^{lc} (1, x_1, x_2, x_3)^T = (1, y_1, y_2, y_3)^T\}$ depicting ellipsoidal surface with semi-axes lengths $\sqrt{\lambda_i/\lambda_0}$, $i = 1, 2, 3$, centered at the origin. We also show that the type-II canonical form Λ_A^{llc} maps the set of all four-vectors $\{(1, x_1, x_2, x_3)^T | x_1^2 + x_2^2 + x_3^2 = 1\}$ on to a spheroidal surface characterised by the four-vectors $(1, y_1, y_2, y_2)^T$ obeying $\frac{y_1^2 + y_2^2}{r_1^2} + \frac{(y_3 - (1-r_0))^2}{r_0^2} = 1$. The geometric intuition underlying the canonical forms is displayed in the figure.

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